## Structural vibration Learning summary

By the end of this chapter you should have learnt about:

- Natural frequencies and mode shapes
- Response of damped single-degree-of-freedom systems
- Response of damped multi-degree-of-freedom systems
- Experimental modal analysis
- Approximate methods
- Vibration control techniques.

# 6.2 Natural frequencies and mode shapes – key points

By the end of this section you should have learnt:

 a systematic approach for setting up the equations of motion for single-degree-of-freedom and lumped mass—spring systems using the following three steps:

- Step 1 convert the physical structure into a dynamic mass-spring model

- Step 2 draw free-body diagram(s)
- Step 3 apply the appropriate form(s) of Newton's second law of motion to give the equation(s) of motion for the system

# 6.2 Natural frequencies and mode shapes – key points

- how to obtain the natural frequency of single-degreeof-freedom systems from the coefficients in the equation of motion
- the equations of motion for lumped mass-spring systems and how to obtain the natural frequencies and the corresponding mode shapes
- how to set up the boundary condition equations for shaft/beam vibration problems and obtain the frequency equation and an expression for the mode shapes.

## 6.3 Response of damped singledegree-of-freedom systems – key points

By the end of this section you should have learnt:

 how to obtain the response of single-degree-offreedom mass—spring-damper systems for the cases of 'free' vibration and for harmonic (sinusoidal) and arbitrary periodic excitation.

## 6.4 Response of damped multi-degreeof-freedom systems – key points

By the end of this section you should have learnt:

- how to uncouple the equations of motion of a proportionally damped multi-degree-of-freedom structure by making use of the orthogonality properties of the modal matrix
- how to formulate response expressions, including the frequency response function.

## 6.5 Experimental modal analysis – key points

By the end of this section you should have learnt:

 how the expression for the frequency response function for a multi-degree-of-freedom structure can be used to identify the natural frequencies, damping and mode shapes of structures from experimental tests.

### 6.6 Approximate methods – key points

By the end of this section you should have learnt:

- how to use Dunkerley's method and Rayleigh's method to obtain estimates of the lowest natural frequency of structures
- how to create a single-degree-of-freedom approximation of a more complex structure and use it to estimate the structure's response.

# 6.7 Vibration control techniques – key points

By the end of this section you should have learnt:

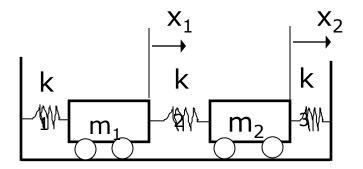
- how to derive the transmissibility expression that gives the ratio of transmitted force to applied force (or of transmitted displacement to applied displacement) as a function of excitation frequency
- how to select suitable isolators to achieve a given isolation efficiency
- how tuned vibration absorbers can suppress resonant vibration in situations where sinusoidal excitation coincides with a natural frequency of the structure.

## Multi-degree of Freedom Systems

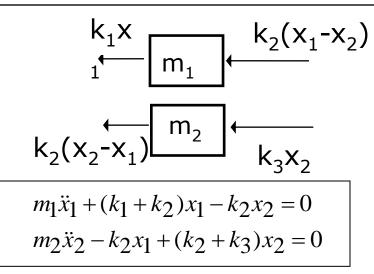
- Motivation: Many systems are too complex to be represented by a single degree of freedom model.
- Objective of this chapter: Understand vibration of systems with more than one degree of freedom.

## Free-vibration of undamped two-degree of freedom system

• We learn how to analyze free vibration by considering an example



### Deriving equations of motion



#### **Interpretation of coefficients**

*First equation*:  $k_1+k_2$  is the force on  $m_1$  needed to move it slowly by one unit while  $m_2$  is held stationary.  $-k_2$  is the force on  $m_1$  need to hold it steady if  $m_2$  is displaced slowly by one unit.

Second equation:  $k_1+k_2$  is the force on  $m_2$  needed to move it slowly by one unit while  $m_1$  is held stationary.  $-k_2$  is the force on  $m_2$  need to hold is steady if  $m_1$  is displaced slowly by one unit.

*Maxwell reciprocity theorem*: the force on  $m_1$  need to hold it steady if  $m_2$  is displaced slowly by one unit=force on  $m_2$  need to hold is steady if  $m_1$  is displaced slowly by one unit.

### Special case

Let m1=m2=m and k1=k2=k3=k. Then:

$$m\ddot{x}_{1} + 2kx_{1} - kx_{2} = 0$$
$$m\ddot{x}_{2} - kx_{1} + 2kx_{2} = 0$$

#### Solution of equations of motion We know from experience that:

 $x_1(t) = A\sin(\omega t)$  $x_2(t) = B\sin(\omega t)$ 

Substituting the above equation to eq. of motion, we obtain two eqs w.r.t. A, B :

 $(-mA\omega^{2} + 2kA - kB)\sin\omega t = 0$  $(-mB\omega^{2} + 2kB - kA)\sin\omega t = 0$ 

The above two equations are satisfied for every t

$$-mA\omega^{2} + 2kA - kB = 0$$
$$-mB\omega^{2} + 2kB - kA = 0$$

## Trivial solution A, B=0. In order to have a nontrivial solution:

$$\det \begin{bmatrix} 2k - m\omega^2 & -k \\ -k & 2k - m\omega^2 \end{bmatrix} = 0$$

$$\omega = \pm \sqrt{\frac{k}{m}}, \pm \sqrt{\frac{3k}{m}}$$

First natural frequency:

$$\omega = \sqrt{\frac{k}{m}} \Longrightarrow A = B$$

The displacement for the first natural frequen  $\begin{pmatrix} A \\ Y_A \end{pmatrix}$ s: This vector is called mode shape. Constant cannot be detern Usually assume first entry is 1. Therefore, mode shape is: Thus, both masses move in phase and the have the same amplitudes.

Second natural frequency:

$$\omega = \sqrt{\frac{3k}{m}} \Longrightarrow A = -B$$

Mode shape for above frequency  $\begin{bmatrix} A \\ -A \end{bmatrix}$ 

Usually assume first entry is 1. Therefore, mode shape is:

 $\begin{pmatrix} l \\ 1 \end{pmatrix}$ 

Displacements of two masses are sums of displacements in the two modes:

$$\binom{x_1(t)}{x_2(t)} = c_1 \binom{1}{1} \sin\left(\sqrt{\frac{k}{m}}t + \psi_1\right) + c_2 \binom{1}{-1} \sin\left(\sqrt{\frac{3k}{m}}t + \psi_2\right)$$

# General expression for vibration of the two-degree of freedom system

$$\begin{pmatrix} x_1(t) \\ x_2(t) \end{pmatrix} = c_1 \begin{pmatrix} A_1 \\ B_1 \end{pmatrix} \sin(\omega_1 t + \psi_1) + c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2)$$

## **Observations:**

System motion is superposition of two harmonic (sinusoidal) motions with frequencies  $\omega_1$  and  $\omega_2$ .

- Participation of each mode depends on initial conditions.
- Four unknowns ( $c_1$ ,  $c_2$ ,  $\psi_1$ ,  $\psi_2$ ) can be found using four initial conditions.
- Can find specific initial conditions so that only one mode is excited.

#### How to solve a free vibration problem involving a two degree of freedom system

1) Write equations of motion for free vibration (no external force or moment) 2) Assume displacements are sinusoidal waves, and plug in equations of motion: Obtain equation:  $[C(\omega)]^* \begin{cases} A \\ B \end{cases} = \begin{cases} 0 \\ 0 \end{cases}$ 

where  $[C(\omega)] = [K] - \omega^2[M]$ 

3) Solve det[C( $\omega$ )]=0, obtain two natural frequencies,  $\omega_1$ ,  $\omega_2$ .

4) Solve,

$$[\mathbf{C}(\boldsymbol{\omega}_1)] * \begin{cases} \mathbf{A} \\ \mathbf{B} \end{cases} = \begin{cases} \mathbf{0} \\ \mathbf{0} \end{cases}$$

Assume A=1, and obtain first mode shape. Obtain second mode shape in a similar manner.

Free vibration response is:

$$\begin{pmatrix} x_1(t) \\ x_2(t) \end{pmatrix} = c_1 \begin{pmatrix} A_1 \\ B_1 \end{pmatrix} \sin(\omega_1 t + \psi_1) + c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) - c_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) + c_2 \begin{pmatrix} A_2 \\ B_2 \end{pmatrix} \sin(\omega_2 t + \psi_2) + c_2$$

5) Find constants from initial conditions.